gaussian-shaped transverse traveling wave is "flipped over" on the *y*-axis, *i*.*e*. *y*o changes to  $-y_0$ . This is equivalent to a phase change, or phase shift of the reflected wave relative to the incident wave of 180°. Thus, the gaussian-shaped transverse traveling wave reflected from a fixed end located at  $x = L$  is of the form  $y(x,t) = -y_0 \exp\{-[(x-2L)+v_x t]^2\}$ , after reflection.

 Another equivalent way to think about this process is to imagine a special kind of mirror at the fixed endpoint  $x = L$ . As the right-moving gaussian-shaped transverse traveling wave,  $y(x,t) = y_0 \exp{-\frac{(x-v_x t)^2}{2}}$  approaches the fixed end, an observer can see this wave directly, but looking in the special mirror, this observer *also* sees a *left*-moving, but *polarity-inverted* gaussian-shaped transverse traveling wave,  $y(x,t) = -y_0$  $\exp\{-[(x-2L)+v_xt)^2\}$  approaching the endpoint at  $x = L$ , but from *behind* the mirror. This left-moving, polarity inverted pulse, approaching from *behind* the end point is known as an *image pulse*. This way of thinking about the problem arises from use of the principle of linear superposition. The fixed-end boundary condition,  $y_{tot}(x=L,t) = 0$  for the incident, right-moving transverse traveling wave and its image pulse, independent of the detailed shape (transverse profiles) of these transverse traveling waves, is generically given by:

$$
y_{\text{tot}}(x=L,t) = f(x-v_x t)|_{x=L} + f(x+v_x t)|_{x=L} = 0
$$

where " $|x=L$ " means evaluating these functions at  $x = L$ . Thus, the fixed end boundary condition,  $y_{\text{tot}}(x=L,t) = 0$  can only be satisfied if  $f(x-v_x t)|_{x=L} = -f(x+v_x t)|_{x=L}$ .

 Note that the peak of the right-moving gaussian-shaped transverse traveling wave arrives at  $x = L$  when  $t = L/v_x$ . The peak location, x of the right-moving gaussian pulse occurs at a given time, *t* when  $(x-v_x t) = 0$ . Thus, at  $t = 0$ ,  $x = 0$ . Note also that the leftmoving gaussian-shaped transverse traveling wave has had its origin shifted, such that when  $t = 0$ , its gaussian peak occurs at  $x = 2L$ , and when  $t = L/v_x$ , its gaussian peak also is at  $x = L$ . The peak, x of this left-moving gaussian pulse occurs at a given time, t when  $[(x-2L)+v_xt] = 0$ . Thus, at  $t = 0$ ,  $x = 2L$  for the left-moving gaussian pulse.

 As time progresses, the two pulses move closer and closer to each other, *passing through* each other at  $x = L$ , and then keep on going. During the time that the two pulses overlap each other, the overall transverse displacement of the string at  $x = L$  is just the *linear superposition* of the two waves, *i*.*e*.

$$
y_{\text{tot}}(x=L, t) = y_{\text{right}}(x=L, t) + y_{\text{left}}(x=L, t) = y_0 \exp\{-(L-v_x t)^2\} - y_0 \exp\{-(L+v_x t)^2\}
$$

In general, for *any* point *x*, and time *t*, we will have, for reflection at a fixed end:

$$
y_{\text{tot}}(x, t) = y_{\text{right}}(x, t) + y_{\text{left}}(x, t) = y_0 \exp\{-(x - v_x t)^2\} - y_0 \exp\{-(x - 2L) + v_x t\}^2
$$

Of course, what we see physically is only the portion of the string for the two waves with  $x \leq L$ . Mathematically, we can represent the wave that is reflected from a fixed end as an image wave, traveling in the opposite direction, and with inverted polarity, to the incoming traveling wave incident on the fixed end. Thus, because it takes a finite time for such a wave to undergo reflection from a fixed end, we can represent this process as a linear superposition of two waves, the original, right-moving transverse traveling wave,