The use of a dual-channel *LIA* for phase-sensitive measurements of periodic complex pressure  $\tilde{p}(\vec{r},t)$  and/or particle velocity  $\vec{u}(\vec{r},t)$  signals associated with a complex sound field  $\tilde{S}(\vec{r},t;\omega)$  has associated with it one other detail which is not immediately apparent in the above formulae.

Consider the use of a dual-channel *LIA* in a typical <u>*phase-sensitive*</u> acoustical physics experiment, such as the typical one shown in the figure below, in which a sound source is excited by a pure-tone/single-frequency sine-wave signal output from a function generator, whose instantaneous output voltage is of the form  $V^{FG}(\omega, t) = V_o^{FG} \cos(\omega t)$ . The sine-wave signal is also used as the <u>reference</u> for the dual-channel *LIA*.

For simplicity's sake, let us imagine that we have an <u>ideal</u> sound source, in that it does <u>not</u> introduce a phase shift of any kind in the process of generating a monochromatic traveling plane wave, which propagates as a <u>free-field</u>. The <u>time-domain</u> representation of the complex overpressure amplitude associated with the <u>free-field</u> monochromatic traveling plane wave propagating in the +ve  $\hat{x}$ -direction is of the form  $\tilde{p}(x,t;\omega) = p_o e^{i(\omega t - kx)} = p_o e^{-ikx} \cdot e^{i\omega t} = \tilde{p}(x,\omega) \cdot e^{i\omega t}$ , where  $\tilde{p}(x,\omega)$  is the <u>frequency-domain</u> representation of the complex over-pressure amplitude

 $p(x, \omega)$  is the <u>frequency-domain</u> representation of the complex over-pressure amplitude associated with the monochromatic traveling plane wave.



If the pressure mic is located at x = 0, then:  $\tilde{p}(x = 0, t) = p_o \cdot e^{i\omega t}$ . The <u>real/in-phase</u> and <u>imaginary/90<sup>o</sup>-out-of-phase/quadrature</u> RMS voltage amplitude components output from the LIA will be  $V_{o_{RMS}}^{Sig}$  and 0, respectively -i.e. the phase  $\varphi_S|_{x=0} = 0$ .

If the *p*-mic sensitivity is  $S_{p\text{-mic}}(mV/Pascal)$ , then the <u>frequency-domain real/in-phase</u> and <u>imaginary/90°-out-of-phase/quadrature</u> RMS components of the complex pressure amplitude  $\tilde{p}(x=0,\omega)$  at x=0 are:  $p_o = V_{o_{RMS}}^{Sig}/S_{p\text{-mic}}(RMS Pascals)$  and: 0(RMS Pascals), respectively.

If the pressure mic is instead moved to x = d, the complex over-pressure amplitude associated with the a *free-field* monochromatic traveling plane wave at x = d is:

$$\tilde{p}(x=d,t;\omega) = p_o e^{i(\omega t - kd)} = p_o e^{-ikd} \cdot e^{i\omega t} = p_o [\cos kd - i\sin kd] \cdot e^{i\omega t} = \tilde{p}(x=d,\omega) \cdot e^{i\omega t}$$

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