

Since $p(\vec{r}, t) = p(r, t) \neq fcn(\theta, \varphi)$, then $\vec{\nabla}p(\vec{r}, t) = \vec{\nabla}p(r, t) = (\partial p(r, t)/\partial r)\hat{r}$ and thus the *instantaneous/physical* (i.e. purely **real time-domain**) particle velocity $\vec{u}(\vec{r}, t)$ associated with a spherically-symmetric monochromatic point sound source can **only** be oriented in the **radial** direction, i.e. $\vec{u}(\vec{r}, t) = u_r(r, t)\hat{r} \neq fcn(\theta, \varphi)$:

$$\frac{\partial \vec{u}(\vec{r}, t)}{\partial t} = -\frac{1}{\rho_o} \vec{\nabla}p(\vec{r}, t) \Rightarrow \frac{\partial u_r(r, t)}{\partial t} \hat{r} = -\frac{1}{\rho_o} \frac{\partial p(r, t)}{\partial r} \hat{r} \Rightarrow \frac{\partial u_r(r, t)}{\partial t} = -\frac{1}{\rho_o} \frac{\partial p(r, t)}{\partial r}$$

$$\text{Now: } \frac{\partial p(r, t)}{\partial r} = \frac{\partial}{\partial r} \left\{ \frac{B_o}{r} \cos(\omega t - kr) \right\} = -\frac{B_o}{r^2} \cos(\omega t - kr) + \frac{kB_o}{r} \sin(\omega t - kr)$$

$$\begin{aligned} \text{Then: } \frac{\partial u_r(r, t)}{\partial t} &= -\frac{1}{\rho_o} \frac{\partial p(r, t)}{\partial r} = -\frac{1}{\rho_o} \left\{ -\frac{B_o}{r^2} \cos(\omega t - kr) + \frac{kB_o}{r} \sin(\omega t - kr) \right\} \\ &= +\frac{B_o}{\rho_o r} \left\{ \frac{1}{r} \cos(\omega t - kr) - k \sin(\omega t - kr) \right\} \end{aligned}$$

Thus:

$$u_r(r, t) = \frac{B_o}{\omega \rho_o r} \left\{ \frac{1}{r} \sin(\omega t - kr) + k \cos(\omega t - kr) \right\} = \frac{B_o k}{\omega \rho_o r} \left\{ \cos(\omega t - kr) + \frac{1}{kr} \sin(\omega t - kr) \right\} \quad (m/s)$$

Using $c = \omega/k$ and $z_o \equiv \rho_o c$ this relation becomes:

$$u_r(r, t) = \frac{B_o}{z_o r} \left\{ \cos(\omega t - kr) + \frac{1}{kr} \sin(\omega t - kr) \right\} \quad (m/s)$$

We can then “complexify” the radial-outgoing spherical over-pressure and particle velocity waves:

$$p(r, t) = \frac{B_o}{r} \cos(\omega t - kr) \Rightarrow \tilde{p}(r, t) = \frac{B_o}{r} \left\{ \cos(\omega t - kr) + i \sin(\omega t - kr) \right\} = \frac{B_o}{r} e^{i(\omega t - kr)}$$

$$\text{and: } u_r(r, t) = \frac{1}{z_o} \frac{B_o}{r} \left\{ \cos(\omega t - kr) + \frac{1}{kr} \sin(\omega t - kr) \right\} \Rightarrow$$

$$\begin{aligned} \tilde{u}_r(r, t) &= \frac{1}{z_o} \frac{B_o}{r} \left[\left\{ \cos(\omega t - kr) + \frac{1}{kr} \sin(\omega t - kr) \right\} + i \left\{ \sin(\omega t - kr) - \frac{1}{kr} \cos(\omega t - kr) \right\} \right] \\ &= \frac{1}{z_o} \frac{B_o}{r} \left[\left\{ \cos(\omega t - kr) + i \sin(\omega t - kr) \right\} - i \frac{1}{kr} \left\{ \cos(\omega t - kr) + i \sin(\omega t - kr) \right\} \right] \\ &= \frac{1}{z_o} \frac{B_o}{r} \left[\left\{ e^{i(\omega t - kr)} \right\} - i \frac{1}{kr} \left\{ e^{i(\omega t - kr)} \right\} \right] = \frac{1}{z_o} \frac{B_o}{r} \left[1 - \frac{i}{kr} \right] e^{i(\omega t - kr)} \left\{ = \frac{1}{z_o} \left[1 - \frac{i}{kr} \right] \tilde{p}(r, t) \right\} \end{aligned}$$