

The real and imaginary components of the ***on-axis frequency-domain*** complex over-pressure are:

$$p_r(r=z, \omega) = \text{Re}\{\tilde{p}(r=z, \omega)\} = p_o \left[\cos kz - \cos kz \sqrt{1 + (a/z)^2} \right]$$

$$p_i(r=z, \omega) = \text{Im}\{\tilde{p}(r=z, \omega)\} = -p_o \left[\sin kz - \sin kz \sqrt{1 + (a/z)^2} \right]$$

The ***magnitude*** of the ***on-axis frequency-domain*** complex over-pressure amplitude is:

$$|\tilde{p}(r=z)| = \sqrt{\tilde{p}(r=z) \tilde{p}^*(r=z)} = \sqrt{2} p_o \left\{ 1 - \cos kz \left[\sqrt{1 + (a/z)^2} - 1 \right] \right\}^{1/2}$$

Using the trigonometric identity $\cos 2A = 1 - 2\sin^2 A$, this can be equivalently written as:

$$|\tilde{p}(r=z)| = p_o \left| \sin \left\{ \frac{1}{2} kz \left[\sqrt{1 + (a/z)^2} - 1 \right] \right\} \right|$$

For $z \gg a$, the Taylor series expansion for $\sqrt{1 + (a/z)^2} \approx 1 + \frac{1}{2}(a/z)^2$ **and** if ***additionally*** $z \gg \frac{1}{2}ka^2 = \frac{1}{2}(2\pi/\lambda)a^2 = \pi a^2/\lambda \equiv \text{Rayleigh length}$, then the ***magnitude*** of the “***far-field***” ***on-axis frequency-domain*** complex over-pressure amplitude is:

$$|\tilde{p}_{far}(r=z)| \approx \frac{1}{2} p_o (a/z) ka = p_o (\pi a^2/\lambda z)$$

Note that the ***magnitude*** of the “***far-field***” ***on-axis frequency-domain*** complex over-pressure is bounded by:

$$0 \leq |\tilde{p}_{far}(r=z)| \leq 2p_o \text{ for } 0 \leq r=z \leq \infty$$

Extrema of $|\tilde{p}(r=z)|$ occur when: $\frac{1}{2} kz \left[\sqrt{1 + (a/z)^2} - 1 \right] = \frac{1}{2} m\pi$, $m = 0, 1, 2, 3, 4, 5, \dots$. Solving the quadratic equation for z , $|\tilde{p}(r=z)|$ extrema occur when: $z_m/a = (a/m\lambda) - (m\lambda/4a)$, $m = 1, 2, 3, 4, 5, \dots$

Coming in from $z = \infty$, the first ***maxima*** $|\tilde{p}(r=z_1)|_{\max}$ occurs at: $z_1/a = (a/\lambda) - (\lambda/4a)$, the next ***maxima*** $|\tilde{p}(r=z_3)|_{\max}$ occurs at: $z_3/a = (a/3\lambda) - (3\lambda/4a)$, and so on. Zeroes of $|\tilde{p}(r=z_m)|$ occur in between the local ***maxima***, when $m = 2, 4, 6, 8, \dots$

Thus, having obtained the ***on-axis*** complex ***time-domain*** over-pressure amplitude $\tilde{p}(r=z, t)$, we next use the {linearized} Euler’s equation for inviscid fluid flow to obtain the ***on-axis*** complex ***time-domain*** particle velocity $\vec{u}(r=z, \omega, t)$:

$$\frac{\partial \vec{u}(r=z, t)}{\partial t} = -\frac{1}{\rho_o} \vec{\nabla} \tilde{p}(r=z, t) = -\frac{1}{\rho_o} \frac{\partial \tilde{p}(r=z, t)}{\partial z} \hat{z}$$