

$$\begin{aligned}
 \vec{\nabla} \tilde{p}(\vec{r}, t) &= \left[ \left\{ \vec{\nabla} |\tilde{p}_o(\vec{r}, \omega)| \right\} \cdot e^{i\varphi_p(\vec{r}, \omega)} + |\tilde{p}_o(\vec{r}, \omega)| \cdot \vec{\nabla} e^{i\varphi_p(\vec{r}, \omega)} \right] \cdot e^{i\omega t} \\
 &= \left[ \frac{\left\{ \vec{\nabla} |\tilde{p}_o(\vec{r}, \omega)| \right\}}{|\tilde{p}_o(\vec{r}, \omega)|} |\tilde{p}_o(\vec{r}, \omega)| \cdot e^{i\varphi_p(\vec{r}, \omega)} + i |\tilde{p}_o(\vec{r}, \omega)| \cdot \left\{ \vec{\nabla} \varphi_p(\vec{r}, \omega) \right\} e^{i\varphi_p(\vec{r}, \omega)} \right] \cdot e^{i\omega t} \\
 &= \left[ \frac{\vec{\nabla} |\tilde{p}_o(\vec{r}, \omega)|}{|\tilde{p}_o(\vec{r}, \omega)|} + i \vec{\nabla} \varphi_p(\vec{r}, \omega) \right] \underbrace{|\tilde{p}_o(\vec{r}, \omega)| \cdot e^{i\varphi_p(\vec{r}, \omega)} \cdot e^{i\omega t}}_{=\tilde{p}(\vec{r}, t)} \\
 &= \left[ \frac{\vec{\nabla} |\tilde{p}_o(\vec{r}, \omega)|}{|\tilde{p}_o(\vec{r}, \omega)|} + i \vec{\nabla} \varphi_p(\vec{r}, \omega) \right] \tilde{p}(\vec{r}, t)
 \end{aligned}$$

The Euler equation for this “generic” 3-D monochromatic traveling wave is:

$$i\omega \cdot \tilde{u}(\vec{r}, t) = -\frac{1}{\rho_o} \left[ \frac{\vec{\nabla} |\tilde{p}_o(\vec{r}, \omega)|}{|\tilde{p}_o(\vec{r}, \omega)|} + i \vec{\nabla} \varphi_p(\vec{r}, \omega) \right] \tilde{p}(\vec{r}, t)$$

or:

$$\begin{aligned}
 \tilde{u}(\vec{r}, t) &= -\frac{1}{i\rho_o\omega} \left[ \frac{\vec{\nabla} |\tilde{p}_o(\vec{r}, \omega)|}{|\tilde{p}_o(\vec{r}, \omega)|} + i \vec{\nabla} \varphi_p(\vec{r}, \omega) \right] \tilde{p}(\vec{r}, t) = +\frac{i}{\rho_o\omega} \left[ \frac{\vec{\nabla} |\tilde{p}_o(\vec{r}, \omega)|}{|\tilde{p}_o(\vec{r}, \omega)|} + i \vec{\nabla} \varphi_p(\vec{r}, \omega) \right] \tilde{p}(\vec{r}, t) \\
 &= +\frac{1}{\rho_o\omega} \left[ i \frac{\vec{\nabla} |\tilde{p}_o(\vec{r}, \omega)|}{|\tilde{p}_o(\vec{r}, \omega)|} - \vec{\nabla} \varphi_p(\vec{r}, \omega) \right] \tilde{p}(\vec{r}, t) = -\frac{1}{\rho_o\omega} \left[ \vec{\nabla} \varphi_p(\vec{r}, \omega) - i \frac{\vec{\nabla} |\tilde{p}_o(\vec{r}, \omega)|}{|\tilde{p}_o(\vec{r}, \omega)|} \right] \tilde{p}(\vec{r}, t)
 \end{aligned}$$

Thus, for a “generic” 3-D monochromatic traveling wave, the complex **time-domain** 3-D particle velocity  $\tilde{u}(\vec{r}, t)$  is related to the complex **time-domain** over-pressure amplitude  $\tilde{p}(\vec{r}, t)$  via the {linearized} Euler equation relation:

$$\boxed{\tilde{u}(\vec{r}, t) = -\frac{1}{\rho_o\omega} \left[ \vec{\nabla} \varphi_p(\vec{r}, \omega) - i \frac{\vec{\nabla} |\tilde{p}_o(\vec{r}, \omega)|}{|\tilde{p}_o(\vec{r}, \omega)|} \right] \tilde{p}(\vec{r}, t)}$$

There are two different kinds of terms/contributions on the RHS of this equation. The first term,  $-\vec{\nabla} \varphi_p(\vec{r}, \omega)$  is the {negative of the} spatial gradient of the **phase** of the complex over-pressure amplitude – note that for this contribution,  $\tilde{u}(\vec{r}, t)$  is **in-phase** with  $\tilde{p}(\vec{r}, t)$ . The second term,  $+i \vec{\nabla} |\tilde{p}_o(\vec{r}, \omega)| / |\tilde{p}_o(\vec{r}, \omega)|$  is the {normalized/fractional} spatial gradient of the complex over-pressure **amplitude** – note that for this contribution,  $\tilde{u}(\vec{r}, t)$  is **90°-out-of-phase** with  $\tilde{p}(\vec{r}, t)$ . Then *e.g.* for the specific case of a monochromatic 3-D traveling plane wave,  $\varphi_p(\vec{r}, \omega) = -\vec{k} \cdot \vec{r}$  and  $\tilde{p}_o(\vec{r}, \omega) = p_o \neq fcn(\vec{r}, \omega)$ , thus:  $\vec{\nabla} \varphi_p(\vec{r}, \omega) = -\vec{\nabla}(\vec{k} \cdot \vec{r}) = -\vec{k}$  and:  $\vec{\nabla} |\tilde{p}_o(\vec{r}, \omega)| = 0$ , hence {here}  $\tilde{u}(\vec{r}, t)$  is **in-phase** with  $\tilde{p}(\vec{r}, t)$  and using  $\omega = ck$  we also see that:  $\tilde{u}(\vec{r}, t) = (\tilde{p}(\vec{r}, t) / \rho_o c) \hat{k}$ .