

Physically, just as the complex scalar electrical **impedance** \tilde{Z}_e is a measure of an electrical device to **impede** the **flow** of a complex scalar AC electrical current $\tilde{I} = \tilde{\vec{J}}_e \cdot \vec{S}_\perp$ (C/s) when a complex scalar AC voltage \tilde{V} is applied across the terminals of the electrical device, the complex 3-D vector acoustic impedance $\tilde{\vec{Z}}_a(\vec{r}, \omega)$ is a measure of the acoustical medium's ability to **impede** the **flow** of a complex acoustic mass current $\tilde{\vec{I}}_a(\vec{r}, \omega) = \tilde{\vec{J}}_a(\vec{r}, \omega) \cdot \vec{S}_\perp = \rho_o \tilde{\vec{u}}(\vec{r}, \omega) \cdot \vec{S}_\perp$ (kg/s) for a complex over-pressure $\tilde{p}(\vec{r}, \omega)$ at point \vec{r} .

Similarly, just as complex scalar electrical **admittance** $\tilde{Y}_e = 1/\tilde{Z}_e$ is a measure of the **ease** with which an electrical device **admits** the **flow** of a complex scalar AC electrical current \tilde{I}_e when a complex scalar AC voltage \tilde{V} is applied across the terminals of the electrical device, the complex 3-D vector acoustic admittance $\tilde{\vec{Y}}_a(\vec{r}, \omega) = 1/\tilde{\vec{Z}}_a(\vec{r}, \omega)$ is a measure of the **ease** with which an acoustical medium's **admits** the **flow** of a complex scalar acoustic mass current $\tilde{\vec{I}}_a(\vec{r}, \omega) = \tilde{\vec{J}}_a(\vec{r}, \omega) \cdot \vec{S}_\perp = \rho_o \tilde{\vec{u}}(\vec{r}, \omega) \cdot \vec{S}_\perp$ (kg/s) in the presence of a complex over-pressure $\tilde{p}(\vec{r}, \omega)$ at the point \vec{r} .

Another way to gain some physical insight into the nature of complex 3-D vector **specific** acoustic impedance $\tilde{\vec{z}}_a(\vec{r}, \omega) = \tilde{p}(\vec{r}, \omega)/\tilde{\vec{u}}(\vec{r}, \omega)$ and complex 3-D vector **specific** acoustic admittance $\tilde{\vec{y}}_a(\vec{r}, \omega) = \tilde{\vec{u}}(\vec{r}, \omega)/\tilde{p}(\vec{r}, \omega) = 1/\tilde{\vec{z}}_a(\vec{r}, \omega)$ of a medium associated with a harmonic sound field is to imagine a physical situation where we set the {magnitude} of the complex scalar over-pressure $\tilde{p}(\vec{r}, \omega)$ to be a constant/fixed value, e.g. $|\tilde{p}(\vec{r}, \omega)| = 1.0 \text{ Pascal}$.

Then, for a harmonic sound field, if the complex 3-D vector **specific** acoustic impedance $\tilde{\vec{z}}_a(\vec{r}, \omega) = \tilde{p}(\vec{r}, \omega)/\tilde{\vec{u}}(\vec{r}, \omega)$ at the point \vec{r} happens to be very **high**, for a fixed complex scalar over-pressure $\tilde{p}(\vec{r}, \omega)$, this tells us that the complex 3-D vector particle velocity $\tilde{\vec{u}}(\vec{r}, \omega)$ at that point must therefore be very **small**, and hence the corresponding complex 3-D vector acoustic mass current density $\tilde{\vec{J}}_a(\vec{r}, \omega) = \rho_o \tilde{\vec{u}}(\vec{r}, \omega)$ at that point must also be very **small**.

Conversely, if for a harmonic sound field the complex 3-D vector **specific** acoustic impedance $\tilde{\vec{z}}_a(\vec{r}, \omega) = \tilde{p}(\vec{r}, \omega)/\tilde{\vec{u}}(\vec{r}, \omega)$ at the point \vec{r} happens to be very **low**, for a fixed complex scalar over-pressure $\tilde{p}(\vec{r}, \omega)$, this tells us that the complex 3-D vector particle velocity $\tilde{\vec{u}}(\vec{r}, \omega)$ at that point must therefore be very **large**, and hence the corresponding complex 3-D vector acoustic mass current density $\tilde{\vec{J}}_a(\vec{r}, \omega) = \rho_o \tilde{\vec{u}}(\vec{r}, \omega)$ at that point must also be very **large**.