

Complex Specific Acoustic Immittances - Admittance and Impedance of a Medium:

The **medium** (solid, liquid or gas) in which sound waves propagate has associated with it the property of how **easy** (or how **difficult**) it is for sound waves to propagate through that medium – the so-called complex **specific acoustic immittances** – complex **specific acoustic admittance** and/or complex **specific acoustic impedance** (the reciprocal of complex **specific acoustic admittance**) give us such information.

For propagation of 1-D sound waves in a medium, the complex **specific acoustic immittances** – i.e. collectively the complex **specific acoustic admittance** and/or complex **specific acoustic impedance** are both well-defined quantities. They are defined in analogy to the complex form of Ohm's Law ($\tilde{V} = \tilde{I}\tilde{Z}$, $\tilde{I} = \tilde{V}\tilde{Y}$) as used e.g. in electrical circuit theory, since complex over-pressure \tilde{p} is the analog of complex AC voltage \tilde{V} , and particle velocity \tilde{u} is ~ the analog of complex AC electric current \tilde{I}_e {Note that $\tilde{J}_a(\vec{r}, t) \equiv \rho_o \tilde{u}(\vec{r}, t)$ ($kg/s-m^2$) is the complex acoustic **mass** current density}, whereas $\tilde{J}_e \equiv \tilde{I}/\tilde{A}_\perp = n_e q_e \tilde{v}_e = \rho_e \tilde{v}$ ($Amp/m^2 = Coul/s-m^2$) is the complex **electrical** current density}. Note also that both \tilde{J}_e and \tilde{J}_a are 3-D **vector** quantities.

Complex Scalar Electrical Immittances:

$$\text{Complex Electrical Impedance: } \tilde{Z}_e(t; \omega) \equiv \frac{\tilde{V}(t; \omega)}{\tilde{I}_e(t; \omega)} \quad (\text{Ohms} = \text{Volts/Amps})$$

$$\text{Complex Electrical Admittance: } \tilde{Y}_e(t; \omega) \equiv \frac{\tilde{I}_e(t; \omega)}{\tilde{V}(t; \omega)} \quad (\text{Siemens} = \text{Ohms}^{-1} = \text{Amps/Volts})$$

If we write out these relations using complex polar notation: $\tilde{V}(t; \omega) = |\tilde{V}(\omega)| e^{i\varphi_V(\omega)} \cdot e^{i\omega t}$, $\tilde{I}_e(t; \omega) = |\tilde{I}_e(\omega)| e^{i\varphi_I(\omega)} \cdot e^{i\omega t}$, then, noting the cancellation of $e^{i\omega t}$ time dependence factors:

$$\tilde{Z}_e(t; \omega) \equiv \frac{\tilde{V}(t; \omega)}{\tilde{I}_e(t; \omega)} = \frac{|\tilde{V}(\omega)| e^{i\varphi_V(\omega)} \cdot e^{i\omega t}}{|\tilde{I}_e(\omega)| e^{i\varphi_I(\omega)} \cdot e^{i\omega t}} = \frac{|\tilde{V}(\omega)| e^{i\varphi_V(\omega)}}{|\tilde{I}_e(\omega)| e^{i\varphi_I(\omega)}} = \frac{|\tilde{V}(\omega)|}{|\tilde{I}_e(\omega)|} e^{i[\varphi_V(\omega) - \varphi_I(\omega)]} = |\tilde{Z}_e(\omega)| e^{i\varphi_Z(\omega)} = \tilde{Z}_e(\omega)$$

$$\tilde{Y}_e(t; \omega) \equiv \frac{\tilde{I}_e(t; \omega)}{\tilde{V}(t; \omega)} = \frac{|\tilde{I}_e(\omega)| e^{i\varphi_I(\omega)} \cdot e^{i\omega t}}{|\tilde{V}(\omega)| e^{i\varphi_V(\omega)} \cdot e^{i\omega t}} = \frac{|\tilde{I}_e(\omega)| e^{i\varphi_I(\omega)}}{|\tilde{V}(\omega)| e^{i\varphi_V(\omega)}} = \frac{|\tilde{I}_e(\omega)|}{|\tilde{V}(\omega)|} e^{i[\varphi_I(\omega) - \varphi_V(\omega)]} = |\tilde{Y}_e(\omega)| e^{i\varphi_Y(\omega)} = \tilde{Y}_e(\omega)$$

Now: $|\tilde{Z}_e(\omega)| = 1/|\tilde{Y}_e(\omega)|$ or: $|\tilde{Y}_e(\omega)| = 1/|\tilde{Z}_e(\omega)|$, and we see that: $\varphi_Z(\omega) = \varphi_V(\omega) - \varphi_I(\omega) = -\varphi_Y(\omega)$, hence: $\tilde{Y}_e(\omega) = |\tilde{Y}_e(\omega)| e^{i\varphi_Y(\omega)} = \{1/|\tilde{Z}_e(\omega)|\} e^{-i\varphi_Z(\omega)} = 1/\{|\tilde{Z}_e(\omega)| e^{i\varphi_Z(\omega)}\} = 1/\tilde{Z}_e(\omega)$. Thus:

$$\tilde{Z}_e(t; \omega) \equiv \frac{\tilde{V}(t; \omega)}{\tilde{I}_e(t; \omega)} = \frac{\tilde{V}(\omega)}{\tilde{I}_e(\omega)} = \tilde{Z}_e(\omega) = \frac{1}{\tilde{Y}_e(\omega)} \quad \text{and:} \quad \tilde{Y}_e(t; \omega) \equiv \frac{\tilde{I}_e(t; \omega)}{\tilde{V}(t; \omega)} = \frac{\tilde{I}_e(\omega)}{\tilde{V}(\omega)} = \tilde{Y}_e(\omega) = \frac{1}{\tilde{Z}_e(\omega)}$$