Euler's Equation for Inviscid Fluid Flow

Euler's equation for *inviscid* (*i.e. dissipationless*) fluid flow is a <u>special/limiting case</u> of the more general {non-linear} Navier-Stokes equation – which expresses Newton's 2nd law of motion for {compressible} fluid flow. The N-S eq'n, in the absence of external driving forces is:

$$\tilde{\rho}(\vec{r},t)\frac{D\tilde{\vec{u}}(\vec{r},t)}{Dt} = -\vec{\nabla}\tilde{p}(\vec{r},t) + \left(\frac{4}{3}\eta + \eta_B\right)\vec{\nabla}\left(\vec{\nabla}\cdot\vec{\vec{u}}(\vec{r},t)\right) - \eta\left(\vec{\nabla}\times\left(\vec{\nabla}\times\vec{\vec{u}}(\vec{r},t)\right)\right)$$

The two <u>dissipative</u> terms on the right-hand side of the Navier-Stokes equation – a non-zero gradient of the divergence of the particle velocity $\vec{\nabla}(\vec{\nabla} \cdot \vec{u}(\vec{r},t))$ and the curl of the <u>vorticity</u> of the particle velocity $\vec{\nabla} \times (\vec{\nabla} \times \vec{u}(\vec{r},t))$ are associated with the <u>coefficient of shear viscosity</u> of the fluid η , and the <u>coefficient of bulk viscosity</u> of the fluid η_B , both of which have SI units of Pascal-seconds (*Pa-s*).

The time derivative term on the left-hand side of the Navier-Stokes equation, $\frac{D\tilde{u}(\vec{r},t)}{Dt}$ is the complex particle <u>acceleration</u> associated with an infinitesimal volume element V of fluid {e.g. air} centered on the space-time point (\vec{r},t) . From dimensional analysis, note that

$$\tilde{\rho}(\vec{r},t)\frac{D\tilde{\vec{u}}(\vec{r},t)}{Dt}\left(\frac{kg-m/s^2}{m^3} = \frac{N}{m^3}\right) \text{ is a force } \underline{density}. \text{ The term } \frac{D}{Dt} \equiv \frac{\partial}{\partial t} + \tilde{\vec{u}}(\vec{r},t)\cdot\vec{\nabla} \text{ is known as}$$

the <u>convective</u> (or <u>substantive</u>, $aka \underline{material}$) derivative, computed from a <u>stationary</u> observer's reference frame, *e.g.* fixed in the <u>laboratory</u>:

$$\frac{D}{Dt} = \frac{\partial}{\partial t} + \frac{\partial \tilde{x}(\vec{r},t)}{\partial t} \frac{\partial}{\partial x} + \frac{\partial \tilde{y}(\vec{r},t)}{\partial t} \frac{\partial}{\partial y} + \frac{\partial \tilde{z}(\vec{r},t)}{\partial t} \frac{\partial}{\partial z}$$
$$= \frac{\partial}{\partial t} + \tilde{u}_{x}(\vec{r},t) \frac{\partial}{\partial x} + \tilde{u}_{y}(\vec{r},t) \frac{\partial}{\partial y} + \tilde{u}_{z}(\vec{r},t) \frac{\partial}{\partial z} = \frac{\partial}{\partial t} + \left(\vec{u}(\vec{r},t)\cdot\vec{\nabla}\right)$$

Euler's equation for inviscid fluid flow is a first-order, linear, homogeneous differential equation, arising from consideration of momentum conservation in a <u>lossless/dissipationless</u> compressible fluid (liquid or gas), that in the absence of external driving forces describes the relationship between complex pressure $\tilde{p}(\vec{r},t)$ and complex particle velocity $\vec{u}(\vec{r},t)$ in the compressible fluid, of volume mass density $\tilde{\rho}(\vec{r},t) (kg/m^3)$. Euler's equation for inviscid fluid flow is thus valid for fluids where the <u>viscosity</u> of the fluid and/or the <u>conduction</u> of <u>heat</u> in the fluid are <u>both</u> zero {or can both be <u>approximated</u> as being <u>negligible</u>}:

$$\tilde{\rho}(\vec{r},t)\frac{D\tilde{\vec{u}}(\vec{r},t)}{Dt} = \tilde{\rho}(\vec{r},t)\left(\frac{\partial\tilde{\vec{u}}(\vec{r},t)}{\partial t} + \left(\tilde{\vec{u}}(\vec{r},t)\cdot\vec{\nabla}\right)\vec{\vec{u}}(\vec{r},t)\right) = -\vec{\nabla}\tilde{p}(\vec{r},t)$$

-1-©Professor Steven Errede, Department of Physics, University of Illinois at Urbana-Champaign, Illinois 2002 - 2017. All rights reserved.